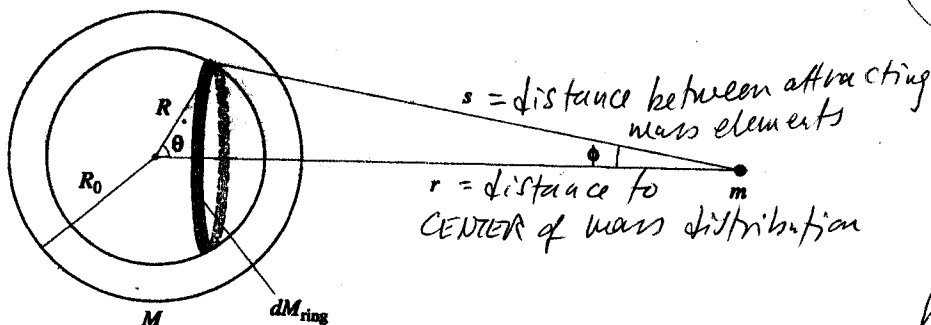
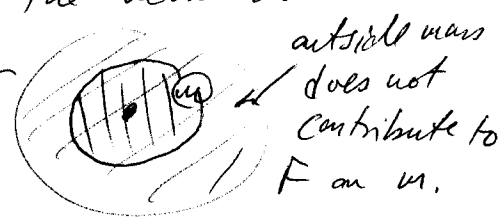


Derive the shell theorem - following Carroll & Ostlie Ch. 2

ZITA
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- The gravitational force due to a spherically symmetric mass M is just the same as that due to a POINT mass M at its center (acting on objects outside the mass distribution) $F = \frac{GmM}{r^2}$
- Corollary: the force on a test mass (m) INSIDE the mass distribution is just the same as that due to ONLY the mass inside the (m) radius, concentrated at the center



$$dF = \frac{Gm dM}{s^2}$$

but since the components of dF due to pairs of mass elements on the ring CANCEL, only the components survive.

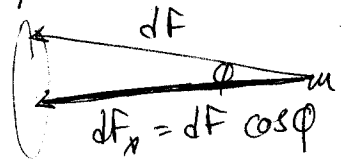
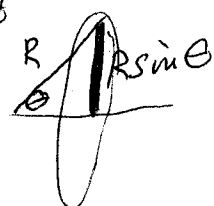
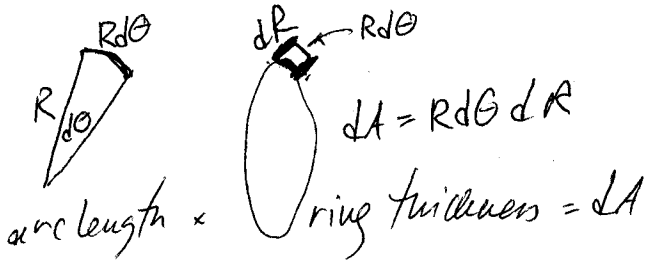


Figure 2.8 The gravitational effect of a spherically symmetric mass distribution.

If dM_{ring} is the mass of the ring being considered, the force exerted by that ring on m is given by

$$dF_{ring} = G \frac{m dM_{ring}}{s^2} \cos \phi.$$

The volume of the ring is its cross-sectional area dA \times its circumference $2\pi R \sin \theta$



$$\begin{aligned} \text{Volume } dV_{ring} &= (dA) 2\pi R \sin \theta \\ &= (R d\theta dR) 2\pi R \sin \theta \end{aligned}$$

Assuming that the mass density, $\rho(R)$, of the extended object is a function of radius only and that the volume of the ring of thickness dR is dV_{ring} , then

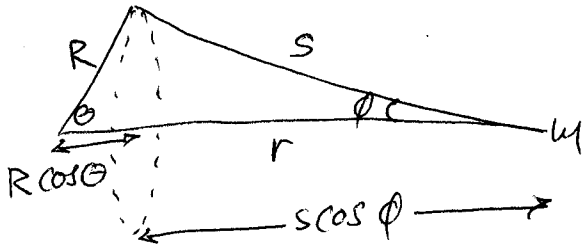
$$\begin{aligned} \text{density} &= \frac{\text{mass}}{\text{volume}} \\ \rho &= \frac{dM}{dV} \end{aligned}$$

$$\begin{aligned} dM_{ring} &= \rho(R) dV_{ring} \\ &= \rho(R) 2\pi R \sin \theta R d\theta dR \\ &= 2\pi R^2 \rho(R) \sin \theta dR d\theta. \end{aligned}$$

So far, we have the force element due to a ring of mass dM on a test mass m :

$$dF_{\text{ring}} = \frac{Gm(dM_{\text{ring}})}{s^2} \cos \phi$$

$$= \frac{Gm}{s^2} [\cos \phi] (2\pi R^2 \rho \sin \theta dR d\theta)$$



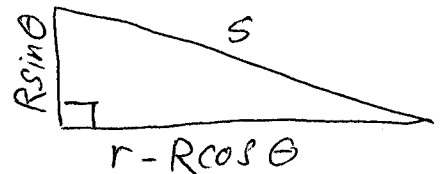
By geometry $\begin{cases} r = R \cos \theta + s \cos \phi, \text{ so} \\ \cos \phi = \frac{r - R \cos \theta}{s} \end{cases}$

where s may be found by the Pythagorean theorem,

$$s = \sqrt{(r - R \cos \theta)^2 + R^2 \sin^2 \theta}$$

$$s = \sqrt{r^2 - 2rR \cos \theta + R^2} = (r^2 - 2rR \cos \theta + R^2)^{1/2}$$

$$s^3 = (r^2 - 2rR \cos \theta + R^2)^{3/2}$$



Substituting into the expression for dF_{ring} , summing over all rings located at a distance R from the center of the mass M (i.e., integrating over all θ from 0 to π for constant R) and then summing over all resultant shells of radius R from $R = 0$ to $R = R_0$ gives the total force of gravity acting on the small mass

$$dF_{\text{ring}} = \frac{Gm}{s^2} \left[\frac{r - R \cos \theta}{s} \right] 2\pi R^2 \rho \sin \theta dR d\theta$$

$$= 2\pi Gm \frac{(r - R \cos \theta)}{s^3} R^2 \rho \sin \theta dR d\theta$$

$$= 2\pi Gm \frac{(r - R \cos \theta)}{(r^2 - 2rR \cos \theta + R^2)^{3/2}} R^2 \rho \sin \theta dR d\theta$$

$$F = Gm \int_0^{R_0} \int_0^\pi \frac{(r - R \cos \theta) \rho(R) 2\pi R^2 (\sin \theta) d\theta}{s^3} dR = Gm \int_0^{R_0} \frac{(r - R \cos \theta) \rho 2\pi R^2}{u^{3/2}} \left(\frac{du}{2rR} \right) dR$$

$$= 2\pi Gm \int_0^{R_0} \int_0^\pi \frac{r R^2 \rho(R) \sin \theta}{(r^2 + R^2 - 2rR \cos \theta)^{3/2}} d\theta dR = -2\pi Gm \int_0^{R_0} \int_0^\pi \frac{R^3 \rho(R) \sin \theta \cos \theta}{(r^2 + R^2 - 2rR \cos \theta)^{3/2}} d\theta dR$$

$$F = 2\pi Gm \int_0^{R_0} \frac{r R^2 \rho}{u^{3/2}} \left(\frac{du}{2rR} \right) dR = F_1$$

$$- 2\pi Gm \int_0^{R_0} \frac{R^3 \rho \cos \theta}{u^{3/2}} \left(\frac{du}{2rR} \right) dR = -F_2$$

The integrations over θ may be carried out by making the change of variable, $u \equiv s^2 = r^2 + R^2 - 2rR \cos \theta$. Then $\cos \theta = (r^2 + R^2 - u)/2rR$ and $\sin \theta d\theta = du/2rR$.

$$\frac{du}{d\theta} = \frac{d}{d\theta} (r^2 + R^2) - 2rR \frac{d}{d\theta} (\cos \theta)$$

$$\leftarrow du = 2rR \sin \theta d\theta$$

$$\leftarrow \frac{du}{d\theta} = 0 - 2rR (-\sin \theta)$$

$$F = F_1 - F_2 \text{ where } F_1 = \pi G m \iint \rho R \frac{du}{u^{3/2}} dR$$

$$\text{and } F_2 = \pi G m \iint \frac{\rho R^2 (\cos \theta)}{r} \frac{du}{u^{3/2}} dR$$

Let's integrate F_1 first.

$$\int \frac{du}{u^{3/2}} = \int u^{-3/2} du = \frac{u^{-1/2}}{-1/2} = \frac{-2}{\sqrt{u}} \Big|_{\theta=0}^{\theta=\pi} = -2 \left[\frac{1}{\sqrt{r^2 + R^2 - 2rR \cos \pi}} - \frac{1}{\sqrt{r^2 + R^2 - 2rR \cos 0}} \right]$$

$$= +2 \left[\frac{-1}{\sqrt{r^2 + R^2 - 2rR}} + \frac{1}{\sqrt{r^2 + R^2 - 2rR}} \right] = 2 \left[\frac{-1}{\sqrt{(r+R)^2}} + \frac{1}{\sqrt{(r-R)^2}} \right]$$

$$\int u^{-3/2} du = 2 \left[\frac{1}{r-R} - \frac{1}{r+R} \right] = 2 \left[\frac{r+R - (r-R)}{(r+R)(r-R)} \right] = \frac{4R}{(r^2 - R^2)}$$

$$\text{Then } F_1 = \pi G m \int \rho R \left(\frac{4R}{r^2 - R^2} \right) dR = \underline{4\pi G m \int \frac{\rho R^2}{r^2 - R^2} dR}$$

$$\text{Next, } F_2 = \pi G m \iint \frac{\rho R^2 (r^2 + R^2 - u)}{2r^2} \frac{du}{u^{3/2}} dR = \frac{\pi G m}{2r^2} \iint \rho R \frac{(r^2 + R^2 - u) du}{u^{3/2}}$$

$$F_2 = \frac{\pi G m}{2r^2} \int \rho R (I_1 - I_2) dR \text{ where } I_1 = \int \frac{r^2 + R^2}{u^{3/2}} du \text{ and } I_2 = \int \frac{u}{u^{3/2}}$$

$$I_1 = (r^2 + R^2) \int u^{-3/2} du = (r^2 + R^2) \frac{4R}{(r^2 - R^2)} \text{ (done above)}$$

$$I_2 = \int \frac{u}{u^{3/2}} = \int u^{-1/2} du = \frac{u^{1/2}}{1/2} \Big|_{\theta=0}^{\theta=\pi} = 2 \left[\sqrt{r^2 + R^2 - 2rR \cos \pi} - \sqrt{r^2 + R^2 - 2rR \cos 0} \right]$$

$$= 2 \left[(r-R) - (r+R) \right] = 2 \left[-2R \right] = -4R$$

$$(ii) \quad I_1 - I_2 = 4R \left[\frac{(r^2 + R^2)}{(r^2 - R^2)} - 1 \right] = 4R \left[\frac{(r^2 + R^2) - (r^2 - R^2)}{(r^2 - R^2)} \right] = 4R \left[\frac{-2R^2}{(r^2 - R^2)} \right]$$

$$\text{Then } F_2 = \frac{\bar{u} G m}{2 r^2} \int \rho R \left(4R \left[\frac{-2R^2}{r^2 - R^2} \right] \right) dR = \frac{4\pi G m}{2} \int \rho \frac{2R^3}{r^2 - R^2} dR$$

$$F_1 - F_2 = 4\pi G m \int \rho \frac{R^2}{r^2 - R^2} \left(1 - \frac{R^2}{r^2} \right) dR$$

$$\left(1 - \frac{R^2}{r^2} = \frac{r^2 - R^2}{r^2} \right)$$

$$F_1 - F_2 = 4\pi G m \int \rho \frac{R^2}{(r^2 - R^2)} \frac{(r^2 - R^2)}{r^2} dR$$

$$F = \frac{4\pi G m}{r^2} \int \rho R^2 dR \quad \checkmark$$

After the appropriate substitutions and integration over the new variable u , the equation for the force becomes

$$F = \frac{Gm}{r^2} \int_0^{R_0} 4\pi R^2 \rho(R) dR.$$

Notice that the integrand is just the mass of a *shell* of thickness dR , having a volume dV_{shell} , or

$$dM_{\text{shell}} = 4\pi R^2 \rho(R) dR$$

$$= \rho(R) dV_{\text{shell}}.$$

Therefore the integrand gives the force on m due to a spherically symmetric mass shell of mass dM_{shell} as

$$dF_{\text{shell}} = \frac{Gm dM_{\text{shell}}}{r^2}.$$

The shell acts gravitationally as if its mass were located entirely at its center. Finally, integrating over the mass shells, we have that the force exerted on m by an extended, spherically symmetric mass distribution is directed along the line of symmetry between the two objects and is given by

$$F = G \frac{Mm}{r^2},$$

just the equation for the force of gravity between two point masses.