

Thus 10 May 07
week 6

Ch 11 - Dipole Radiation - Do #1, 5

*1. Problem 11.1 Check that the retarded potentials of an oscillating dipole (Eqs. 11.12 and 11.17) satisfy the Lorentz gauge condition. Do not use approximation 3.

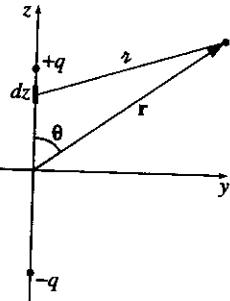
$$\nabla \cdot \mathbf{A} = \frac{\partial V}{\partial t} \text{ for Lorentz Gauge}$$

Putting Eqs. 11.9 and 11.11 into Eq. 11.5, we obtain the potential of an oscillating perfect dipole:

$$V(r, \theta, t) = \frac{\rho_0 \cos \theta}{4\pi \epsilon_0 r} \left\{ -\frac{\omega}{c} \sin[\omega(t - r/c)] + \frac{1}{r} \cos[\omega(t - r/c)] \right\}. \quad (11.12)$$

$$\mathbf{A}(r, \theta, t) = -\frac{\mu_0 \rho_0 \omega}{4\pi r} \sin[\omega(t - r/c)] \hat{z}$$

Small off LF



V is in spherical coords, \mathbf{A} is in cartesian

$$\hat{z} = \cos \theta \hat{i} - \sin \theta \hat{o}$$

$$\begin{aligned} \frac{\partial V}{\partial t} : \quad & \frac{\partial}{\partial t} \sin\left[\omega t - \frac{\omega r}{c}\right] = \omega \cos\left(\omega t - \frac{\omega r}{c}\right) \\ & \frac{\partial}{\partial t} \cos\left[\omega t - \frac{\omega r}{c}\right] = -\omega \sin\left(\omega t - \frac{\omega r}{c}\right) \end{aligned}$$

$$\frac{\partial V}{\partial t} = \frac{\rho_0 \cos \theta}{4\pi \epsilon_0 r} \left\{ -\frac{\omega^2}{c} \cos\left(\omega t - \frac{\omega r}{c}\right) - \frac{\omega}{r} \sin\left(\omega t - \frac{\omega r}{c}\right) \right\} \quad K = \frac{\mu_0 \rho_0 \omega}{4\pi}$$

$$\mathbf{A} = \frac{K}{r} \sin\left(\omega t - \frac{\omega r}{c}\right) [\cos \theta \hat{i} - \sin \theta \hat{o}] \quad \text{in spherical}$$

$$\begin{aligned} \nabla \cdot \mathbf{A} &= \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 \frac{K}{r} \sin\left(\omega t - \frac{\omega r}{c}\right) \cos \theta \right] = \frac{K \cos \theta}{r^2} \frac{\partial}{\partial r} [r \sin\left(\omega t - \frac{\omega r}{c}\right)] \\ &= \frac{K \cos \theta}{r^2} \left[\sin \omega t + \frac{\omega}{c} \cos \omega t \right] = \frac{K}{r^2} \cos \theta \left[\sin \omega t - \frac{\omega r}{c} \cos \omega t \right] \end{aligned}$$

$$\begin{aligned} \nabla \cdot \mathbf{A}_0 &= \frac{1}{r^2} \frac{\partial}{\partial \theta} \left[\sin \theta \frac{K}{r} \sin\left(\omega t - \frac{\omega r}{c}\right) (-\sin \theta) \right] \\ &= -\frac{K \sin(\omega t - \frac{\omega r}{c})}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \sin^2 \theta = -\frac{K \sin \omega t}{r^2 \sin \theta} 2 \sin \theta \cos \theta \end{aligned}$$

$$(\nabla \cdot A)_0 = -\frac{2k}{r^2} \sin \omega t \cos \theta$$

$$(\nabla \cdot A)_r = \frac{k}{r^2} \cos \theta \left[\sin \omega t - \frac{wr}{c} \cos \omega t \right]$$

$$\nabla \cdot A = \frac{k}{r^2} \cos \theta \left[-\sin \left(\omega t - \frac{wr}{c} \right) - \frac{wr}{c} \cos \left(\omega t - \frac{wr}{c} \right) \right]$$

$$= \frac{\mu_0 \rho_0 w}{4\pi r^2} \cos \theta \left[\sin \left(\omega t - \frac{wr}{c} \right) + \frac{wr}{c} \cos \left(\omega t - \frac{wr}{c} \right) \right]$$

$$-\mu_0 E_0 \frac{\partial V}{\partial t} = -\frac{\mu_0 \epsilon_0 \rho_0 \cos \theta}{4\pi \epsilon_0 r} \left[-\frac{w^2}{c} \cos \left(\omega t - \frac{wr}{c} \right) - \frac{w}{r} \sin \left(\omega t - \frac{wr}{c} \right) \right]$$

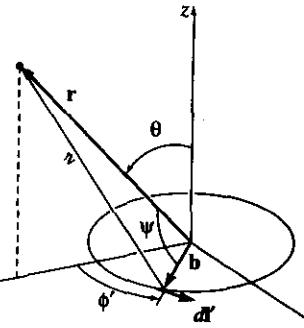
= $\nabla \cdot A$ ✓ ; Lorentz Gauge.

Problem 11.5 Calculate the electric and magnetic fields of an oscillating magnetic dipole without using approximation 3. [Do they look familiar? Compare Prob. 9.33.] Find the Poynting vector, and show that the intensity of the radiation is exactly the same as we got using approximation 3.

vector potential of an oscillating perfect magnetic dipole is

$$\mathbf{A}(r, \theta, t) = \frac{\mu_0 m_0}{4\pi} \left(\frac{\sin \theta}{r} \right) \left\{ \frac{1}{r} \cos[\omega(t - r/c)] - \frac{\omega}{c} \sin[\omega(t - r/c)] \right\} \hat{\phi}. \quad (11.33)$$

$$\begin{aligned} \mathbf{E} &= -\frac{\partial \mathbf{A}}{\partial t} = \frac{\mu_0 m_0}{4\pi c r} \left[\frac{1}{r} \left(\frac{1}{c} \cos \omega t \sin \theta - \frac{\omega}{c} \sin \omega t \cos \theta \right) - \frac{\omega}{c} \left(\frac{1}{c} \right) \sin^2 \frac{\omega}{c} r \right] \hat{\phi} \\ &= \frac{\mu_0 m_0}{4\pi c r} \cos \omega t \sin \theta \left(\frac{1}{c} - \frac{\omega}{c^2} \right) \hat{\phi} \end{aligned}$$



$$\mathbf{B} = \nabla \times \mathbf{A} = \frac{1}{r} \sin \theta \left[\frac{\partial}{\partial \theta} r A_\phi \right] \hat{r} - \left[\frac{1}{r} \frac{\partial}{\partial r} r A_\phi \right] \hat{\theta}$$

$$\begin{aligned} r A_\phi &= \frac{\mu_0 m_0}{4\pi} \sin \theta \left\{ \frac{1}{r} \cos(\omega t - \frac{\omega r}{c}) - \frac{\omega}{c} \sin(\omega t - \frac{\omega r}{c}) \right\} \\ \frac{\partial}{\partial r} r A_\phi &= \frac{\mu_0 m_0}{4\pi} \sin \theta \left\{ -\frac{1}{r^2} \cos \omega t + \frac{1}{r} \left(-\frac{\omega}{c} \right) (-\sin \omega t) + \frac{\omega^2}{c^2} \cos \omega t \right\} \\ &= \frac{\mu_0 m_0}{4\pi} \sin \theta \left\{ \cos \omega t \left(\frac{\omega^2}{c^2} - \frac{1}{r^2} \right) + \sin \omega t \left(\frac{\omega}{rc} \right) \right\} \end{aligned}$$

$$\sin \theta A_\phi = \frac{\mu_0 m_0}{4\pi r} \left[\frac{1}{r} \cos \omega t - \frac{\omega}{c} \sin \omega t \right] \sin^2 \theta$$

$$\frac{\partial}{\partial \theta} \sin \theta A_\phi = \frac{\mu_0 m_0}{4\pi r} \left[\frac{1}{r} \cos \omega t - \frac{\omega}{c} \sin \omega t \right] 2 \sin \theta \cos \theta$$

$$\begin{aligned} \mathbf{B} &= \frac{2 \mu_0 m_0}{4\pi r^2} \left[\frac{1}{r} \cos \omega t - \frac{\omega}{c} \sin \omega t \right] \hat{r} - \frac{\mu_0 m_0}{4\pi r} \sin \theta \left[\cos \omega t \left(\frac{\omega^2}{c^2} - \frac{1}{r^2} \right) \right. \\ &\quad \left. + \sin \omega t \left(\frac{\omega}{rc} \right) \right] \hat{\theta} \\ S &= \frac{1}{\mu_0} \mathbf{E} \cdot \mathbf{B} = \frac{1}{\mu_0} \begin{vmatrix} \hat{r} & \theta & \phi \\ E_r & E_\theta & E_\phi \\ B_r & B_\theta & B_\phi \end{vmatrix} = \frac{1}{\mu_0} \left[\hat{r} \left(-E_\theta B_\phi - \hat{\theta} (-E_\phi B_\theta) \right) \right. \\ &\quad \left. + \sin \theta \left(\frac{\omega}{rc} \right) \right] \end{aligned}$$

Same as Prob 9.33 (p. 412) where we found that

$$\begin{aligned} S &= \frac{\mu_0 m_0^2 \omega^3}{16\pi^2 c^2} \left(\frac{\sin \theta}{r^2} \right) \left\{ \frac{2 \cos \theta}{r} \left[\left(1 - \frac{c^2}{\omega^2 r^2} \right) \sin \omega t \cos \theta + \frac{\omega}{\omega r} (\cos^2 \theta - \sin^2 \theta) \right] \hat{\theta} \right. \\ &\quad \left. + \sin \left[\left(-\frac{\omega}{r} + \frac{c^2}{\omega^2 r^2} \right) \sin \omega t \cos \theta + \frac{\omega}{c} \cos^2 \theta + \frac{c}{\omega r^2} (\sin^2 \theta + \cos^2 \theta) \right] \hat{r} \right\} \end{aligned}$$

where $\omega = \omega(1 - \frac{c}{\omega})$. Intensity is $\langle S \rangle = \frac{\mu_0 m_0^2 \omega^4 \sin^2 \theta}{32\pi^2 c^3 r^2} \hat{r}$ same as (11.39)

Problem 11.6 Find the radiation resistance (Prob. 11.3) for the oscillating magnetic dipole
 Fig. 11.8. Express your answer in terms of λ and b , and compare the radiation resistance
 to the electric dipole. [Answer: $3 \times 10^5 (b/\lambda)^4 \Omega$]

$$I^2 R = I_0^2 R^2 \cos^2 \omega t \rightarrow \langle P \rangle = \frac{1}{2} I_0^2 R = \frac{\mu_0 \mu_0^2 \omega^4}{12 \pi c} = \frac{\mu_0 \pi^2 b^4 \lambda^3}{12 \pi c}$$

$$IR = \mu_0 \frac{\pi b^4 \omega^4}{6c^3} . \text{ Since } \omega = \frac{2\pi c}{\lambda},$$

$$R = \frac{\mu_0 \pi^2 b^4}{6c^3} \frac{16\pi^4 c^4}{\lambda^4} = \frac{8}{3} \pi^5 \mu_0 c \left(\frac{b}{\lambda}\right)^4$$

$$= \frac{8}{3} \pi^5 \left(4\pi \times 10^{-7} \frac{C^2}{N} \right) \left(3 \times 10^8 \frac{m}{s} \right)^2 \left(\frac{b}{\lambda}\right)^4 = 3 \times 10^5 \left(\frac{b}{\lambda}\right)^4 R$$

much smaller than for electric dipole.

$$\text{Ex: } \frac{b=5 \text{ cm}}{\lambda=10^3 \text{ m}} \rightarrow R_B = 3 \times 10^5 \left(\frac{5 \times 10^{-5}}{10^3}\right)^4 = 2 \times 10^{-12} \Omega$$

$\sim 10^{-16} R_E$